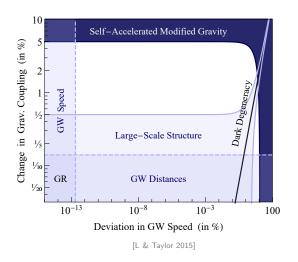
# Challenges to Cosmic Self-Acceleration in Modified Gravity from Gravitational Waves & Large-Scale Structure

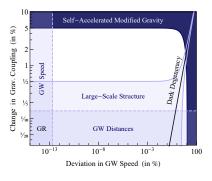
#### Lucas Lombriser

Institute for Astronomy, University of Edinburgh Département de Physique Théorique, Université de Genève

> IPhT, CEA Saclay 26 Sep, 2017

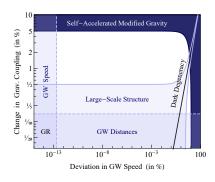






[L & Taylor 2015]

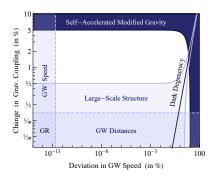
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[L & Taylor 2015]

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(except for usual suspects: quintessence, k-essence)



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- Conclusion 1: LSS without GWs cannot fully discriminate self-accelerated MG from Λ.
- Conclusion 2: Horndeski gravity will no longer be a candidate for explaining cosmic acceleration by October.
   (except for usual suspects: quintessence, k-essence)
- Conclusion 3: Impending death of the Galileon.

## Horndeski scalar-tensor action

Horndeski action:

$$\begin{split} S &= \frac{1}{2\kappa^2} \int d^4x \sqrt{-g} \left\{ G_2(\varphi,X) - G_3(\varphi,X) \Box \varphi \right. \\ &+ \left. G_4(\varphi,X) R + \frac{\partial G_4}{\partial X} \left[ (\Box \varphi)^2 - (\nabla_{\mu} \nabla_{\nu} \varphi)^2 \right] \right. \\ &+ \left. G_5(\varphi,X) G_{\mu\nu} \nabla^{\mu} \nabla^{\nu} \varphi \right. \\ &- \left. \frac{1}{6} \frac{\partial G_5}{\partial X} \left[ (\Box \varphi)^3 - 3 \Box \varphi (\nabla_{\mu} \nabla_{\nu} \varphi)^2 + 2 (\nabla_{\mu} \nabla_{\nu} \varphi)^3 \right] \right. \\ &+ \mathcal{L}_{\mathrm{m}}(g_{\mu\nu},\psi_i) \left. \right\}, \end{split}$$

where 
$$X \equiv -\frac{1}{2}(\partial_{\mu}\varphi)^2$$

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$$\frac{d^2\tilde{a}}{d\tilde{t}^2} = \frac{1}{\sqrt{\Omega}} \left[ \left( 1 + \frac{1}{2} \frac{\Omega'}{\Omega} \right) \frac{d^2 a}{dt^2} + \frac{a H^2}{2} \left( \frac{\Omega'}{\Omega} \right)' \right] \le 0$$



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#### Self-acceleration

The breaking of the strong (or weak) equivalence principle in the cosmological background is responsible for cosmic acceleration.



#### Cosmological background and linear perturbations

H(t): Hubble parameter

Creminelli *et al.* (2008); Park *et al.* (2010); Gubitosi *et al.* (2012); Bloomfield *et al.* (2012); Bellini & Sawicki (2014); Gleyzes *et al.* (2014) . . .

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 $\alpha_{\rm T}(t)$ : Tensor speed alteration  $(c_{\rm T}^2=1+\alpha_{\rm T})$ 

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• Consistency relation:  $\Omega = M^2 c_{\rm T}^2$ , where  $\alpha_{\rm M} = (M^2)'/M^2$ 

[L & Taylor (2015)]



- Modification as an effective fluid  $\rightarrow G^{\mu\nu} \equiv R^{\mu\nu} \frac{1}{2} g^{\mu\nu} R \equiv \kappa^2 \left( T_{\rm m}^{\mu\nu} + T_{\rm eff}^{\mu\nu} \right) = \kappa^2 T^{\mu\nu}$
- $g^{\mu \nu} + \delta g^{\mu \nu} o$ 4 scalar metric perturbations
- $T^{\mu\nu} + \delta T^{\mu\nu} \rightarrow$  4 matter perturbations
- Einstein & conservation equations fix 4
- 2 can be removed from fixing to a gauge
- 2 closure relations needed from modified gravity

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$$ds^2 = -(1 + \Psi)dt^2 + a(t)^2(1 + 2\Phi)dx^2$$

Conservation equations unchanged; modified Einstein equations:

$$k^{2}\Psi = -\frac{\kappa^{2}}{2}\mu(a,k)\bar{\rho}_{m}a^{2}\Delta_{m}$$
  
$$\Phi = -\gamma(a,k)\Psi$$

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Closure relations:

- $\Lambda$ CDM:  $\mu = 1$ ;  $\gamma = 1$
- Horndeski (quasistatic):  $\mu=h_1\left(\frac{1+h_4k^2}{1+h_5k^2}\right)$ ;  $\gamma=h_2\left(\frac{1+h_3k^2}{1+h_4k^2}\right)$



# Linear Shielding

#### Linearly shielded Horndeski scalar-tensor theory

- $\lim_{k\to\infty} \mu(a,k) = \gamma(a,k) = 1$  with 3 free functions of time (1 acting only beyond QS limit)
- $\mu(a, k) = \gamma(a, k) = 1$ with 2 free functions of time
- Can set  $H = H_{\Lambda CDM}$  on top of that

[L & Taylor (2014)]

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# A Dark Degeneracy

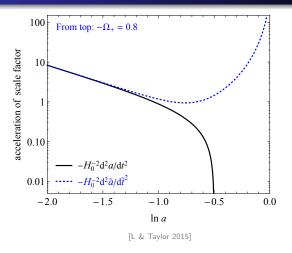
#### A self-accelerated Horndeski MG degenerate with $\Lambda CDM$

- Choose  $\Lambda$ CDM H(t) (in Jordan frame): fixes 1st EFT funct.
- $\Omega_{+} \lesssim -0.1$ : fixes 2nd EFT function
- Apply linear shielding conditions: fixes 3rd & 4th EFT function
- Set  $c_s = c \equiv 1$ : fixes 5th EFT function
- Free of ghost and gradient instabilities

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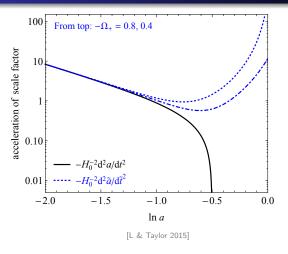
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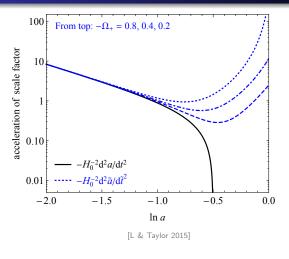
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,  $n = 4$ 





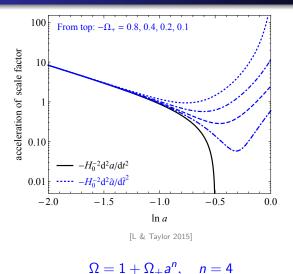
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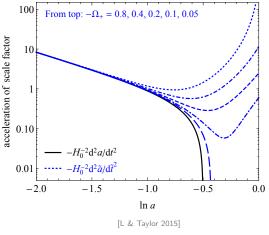


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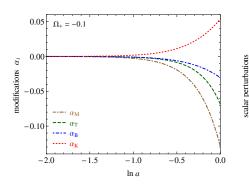
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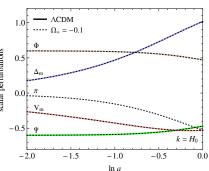
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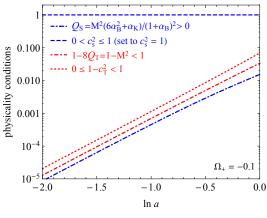
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## Stability & physicality



[L & Taylor 2015]

• Propagation of gravitational waves:

$$h_{ij}^{\prime\prime} + \left(\alpha_{\mathbf{M}} + 3 + \frac{H^{\prime}}{H}\right)h_{ij}^{\prime} + (1 + \alpha_{\mathbf{T}})k_{H}^{2}h_{ij} = 0$$

- Different propagation speed: can be tested by comparing arrival time of signals
- Different damping of GW amplitude: can be tested with Standard Sirens

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#### Requirements for Dark Degeneracy $(H = H_{\Lambda CDM})$

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- Relation between  $\alpha_{\rm M}$  and  $\alpha_{\rm T}$ :

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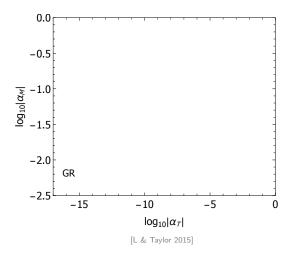
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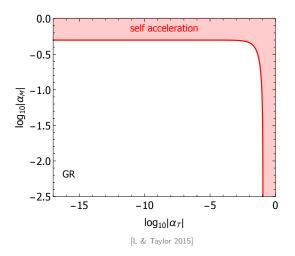
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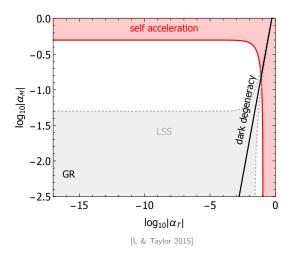
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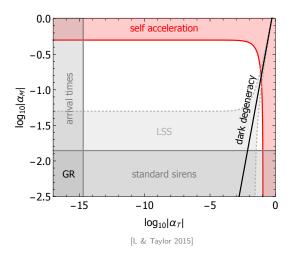
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- Assume  $\alpha_{\rm T} \simeq 0$  ( $c_{\rm T}=1$ ) and  $H=H_{\Lambda {\rm CDM}}$  cosmic rays, binary pulsars, LIGO&VIRGO GW+GRB 17.8.17?
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$$\left| \frac{\Omega'}{\Omega} \right| = \left| \alpha_{\mathrm{M}} + \frac{\alpha_{\mathrm{T}}'}{1 + \alpha_{\mathrm{T}}} \right| \gtrsim \mathcal{O}(1)$$

• Minimal acceleration:

$$\begin{split} \frac{d^2 \tilde{a}}{d \tilde{t}^2} &= \frac{1}{\sqrt{\Omega}} \left[ \left( 1 + \frac{1}{2} \frac{\Omega'}{\Omega} \right) \frac{d^2 a}{d t^2} + \frac{a \, H^2}{2} \left( \frac{\Omega'}{\Omega} \right)' \right] \leq 0 \\ \Rightarrow & \left( 1 + \frac{H'}{H} \right) \left( 1 + \frac{1}{2} \alpha_M \right) + \frac{1}{2} \alpha_M' \leq 0 \\ \Rightarrow & \kappa^2 M^2 \leq \left( \frac{a_{\rm acc}}{a} \right)^2 e^{C(\chi_{\rm acc} - \chi)} \,, \quad C = 2 H_0 a_{\rm acc} \sqrt{3(1 - \Omega_{\rm m})} \end{split}$$

• Minimise modification in growth of structure:  $\alpha_{B} = \alpha_{M}$  follows from

$$\mu_{\infty} = \frac{2(\alpha_{\rm B} - \alpha_{\rm M})^2 + \alpha c_{\rm s}^2}{\alpha c_{\rm s}^2 \kappa^2 M^2}$$

and  $M^2$ ,  $\alpha$ ,  $c_{\rm s}^2 > 0$  for stability

• Minimal self-acceleration:  $\mu = (\kappa^2 M^2)^{-1} \ge 1$  and  $\gamma = 1$  with  $\mu(a \le a_{\rm acc} \simeq 0.6) = 1$  increasing to  $\mu(a = 1) \simeq 1.04$ 

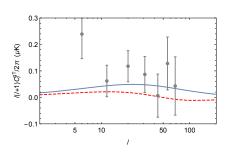
• Minimise modification in growth of structure:  $\alpha_{\rm B}=\alpha_{\rm M}$  follows from

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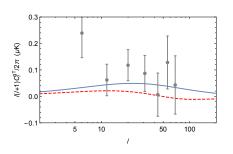
- Background:
   SN Ia, BAO, H<sub>0</sub>, CMB
- Perturbations:
   CMB (Planck 2015), E<sub>G</sub>
   galaxy-ISW
- ISW sensitive to  $\Sigma' = -\alpha_{\rm M} \Sigma$ where  $\Sigma = (1 + \gamma)\mu/2$
- Overall:  $3\sigma$  worse fit than  $\Lambda \mathrm{CDM}$  strong evidence for  $\Lambda$  ( $B \simeq 39$ )





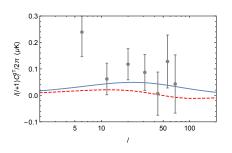
- Background:
   SN Ia, BAO, H<sub>0</sub>, CMB
- Perturbations:
   CMB (Planck 2015), E<sub>G</sub>,
   galaxy-ISW
- ISW sensitive to  $\Sigma' = -\alpha_{\rm M} \Sigma$ where  $\Sigma = (1 + \gamma)\mu/2$
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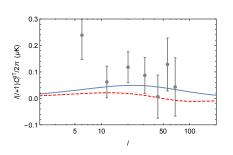
- Background:
   SN Ia, BAO, H<sub>0</sub>, CMB
- Perturbations:
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- Background:
   SN Ia, BAO, H<sub>0</sub>, CMB
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   CMB (Planck 2015), E<sub>G</sub>,
   galaxy-ISW
- ISW sensitive to  $\Sigma' = -\alpha_{\rm M} \Sigma$ where  $\Sigma = (1 + \gamma)\mu/2$
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- Background:
   SN Ia, BAO, H<sub>0</sub>, CMB
- Perturbations: CMB (Planck 2015),  $E_G$ , galaxy-ISW
- ISW sensitive to  $\Sigma' = -\alpha_{\rm M} \Sigma$ where  $\Sigma = (1 + \gamma)\mu/2$
- Overall:  $3\sigma$  worse fit than  $\Lambda {\rm CDM}$  strong evidence for  $\Lambda$   $(B\simeq39)$

[L & Lima (2016)]

Compare: f(R), DGP, Galileon [L, Slosar, Seljak, Hu '10; L, Hu, Fang, Seljak '09; Barreira, Li, Baugh, Pascoli '14]

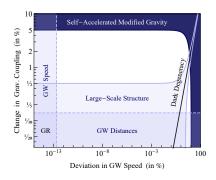


Minimal self-acceleration
Incompatibility with observations

# **NGC 4993**



#### Conclusions



[L & Taylor 2015]

- Conclusion 1: LSS without GWs cannot fully discriminate self-accelerated MG from Λ.
  - → parametrised tests?
- Conclusion 2: Horndeski gravity may no longer be a candidate for explaining cosmic acceleration by October. (except for usual suspects: quintessence, k-essence)
  - $\rightarrow$  WL, gCMB $\phi$ X, Std Sirens
- Conclusion 3: Impending death of the (cov.) Galileon?

Minimal self-acceleration Incompatibility with observations Conclusions/Outlook

# Thank you!

Postdoc & PhD positions at University of Geneva advertised soon

